3.2 GaAs GaAs has an effective density of states at the conduction CB $N_{c}$ of $4.7 \times 10^{17} \mathrm{~cm}^{-3}$ and an effective density of states at the VB edge $N_{v}$ of $7 \times 10^{18} \mathrm{~cm}^{-3}$. Given its bandgap $E_{g}$ of 1.42 eV calculate the intrinsic concentration and the intrinsic resistivity at room temperature (take as 300 K ). Where is the Fermi level? Assuming that $N_{c}$ and $N_{v}$ scale as $T^{3 / 2}$, what would be the intrinsic concentration at $100{ }^{\circ} \mathrm{C}$ ? If this GaAs crystal is doped with $10^{18}$ donors $\mathrm{cm}^{-3}$ (such as Te ), where is the new Fermi level and what is the resistivity of the sample? The drift mobilities in GaAs are shown in Table 3Q2

## Table 3Q2

Dopant impurities scatter carriers and reduce the drift mobility ( $\mu_{e}$ for electrons and $\mu_{h}$ for holes).

| Dopant concentration $\left(\mathrm{cm}^{-3}\right)$ | 0 | $10^{15}$ | $10^{16}$ | $10^{17}$ | $10^{18}$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $\mu_{e}\left(\mathrm{~cm}^{2} \mathrm{~V}^{-1} \mathrm{~s}^{-1}\right)$ | 8500 | 8000 | 7000 | 4000 | 2400 |
| $\mu_{h}\left(\mathrm{~cm}^{2} \mathrm{~V}^{-1} \mathrm{~s}^{-1}\right)$ | 400 | 380 | 310 | 220 | 160 |

## Solution

The intrinsic concentration is

$$
n_{i}=\left(N_{c} N_{v}\right)^{1 / 2} \exp \left(-\frac{E_{g}}{2 k_{B} T}\right)
$$

so that

$$
\begin{gathered}
n_{i}=\left[\left(4.7 \times 10^{17} \mathrm{~cm}^{-3}\right)\left(7 \times 10^{18} \mathrm{~cm}^{-3}\right)\right]^{12} \exp \left[-\frac{1.42 \mathrm{eV}}{2\left(8.6174 \times 10^{-5} \mathrm{eV} \mathrm{~K}^{-1}\right)(300 \mathrm{~K})}\right]= \\
\quad=2.223 \times 10^{12} \mathrm{~m}^{-3}=2.223 \times 10^{6} \mathrm{~cm}^{-3}
\end{gathered}
$$

The conductivity is

$$
\begin{aligned}
\sigma=e n_{i} & \left(\mu_{e}+\mu_{h}\right) \\
& =\left(1.608 \times 10^{-19} \mathrm{C}\right)\left(2.223 \times 10^{6} \mathrm{~cm}^{-3}\right)\left(8500 \mathrm{~cm}^{2} \mathrm{~V}^{-1} \mathrm{~s}^{-1}+400 \mathrm{~cm}^{2} \mathrm{~V}^{-1} \mathrm{~s}^{-1}\right) \\
& =3.17 \times 10^{-9} \Omega^{-1} \mathrm{~cm}^{-1}
\end{aligned}
$$

The resistivity is

$$
\rho=\frac{1}{\sigma}=\frac{1}{\left(3.17 \times 10^{-9} \Omega^{-1} \mathrm{~cm}^{-1}\right)}=3.16 \times 10^{8} \Omega \mathrm{~cm}
$$

The energy distance between the bottom of the conduction band $E_{c}$ and Fermi level of the intrinsic GaAs $E_{F i}$ is

$$
E_{c}-E_{F i}=k_{B} T \ln \left(\frac{N_{c}}{n_{i}}\right)
$$

so that

$$
E_{c}-E_{F i}=\left(8.617 \times 10^{-5} \mathrm{eV} \mathrm{~K}^{-1}\right)(300 \mathrm{~K}) \ln \left[\frac{\left(4.7 \times 10^{17} \mathrm{~cm}^{-3}\right)}{\left(2.223 \times 10^{6} \mathrm{~cm}^{-3}\right)}\right]=0.675 \mathrm{eV}
$$

If $N_{c}$ and $N_{v}$ scale with temperature as $T^{3 / 2}$ then the intrinsic concentration at a given temperature $T$ is given by the expression

$$
n_{i}(T)=\left[N_{c}(300 \mathrm{~K}) N_{v}(300 \mathrm{~K})\right]^{1 / 2}\left(\frac{T}{300 \mathrm{~K}}\right)^{3 / 2} \exp \left(-\frac{E_{g}}{2 k_{B} T}\right),
$$

so at $100^{\circ} \mathrm{C}(373 \mathrm{~K})$ we will have

$$
\begin{aligned}
& n_{i}(373 \mathrm{~K})=\left[\left(4.7 \times 10^{17} \mathrm{~cm}^{-3}\right)\left(7 \times 10^{18} \mathrm{~cm}^{-3}\right)\right]^{12} \\
& \quad \times\left(\frac{373 \mathrm{~K}}{300 \mathrm{~K}}\right)^{3 / 2} \exp \left[-\frac{1.42 \mathrm{eV}}{2\left(8.6174 \times 10^{-5} \mathrm{eVK}^{-1}\right)(373 \mathrm{~K})}\right] \\
& =6.61 \times 10^{8} \mathrm{~cm}^{-3} .
\end{aligned}
$$

For the case when the same GaAs crystal is doped with $10^{18} \mathrm{~cm}^{-3}$ donors, one might be tempted to think that $n=N_{d}$ and to apply the usual expression for the calculation of the Fermi level position:

$$
E_{c}-E_{F n}=k_{B} T \ln \left(\frac{N_{c}}{n}\right)=k_{B} T \ln \left(\frac{N_{c}}{N_{d}}\right)
$$

The result from such calculation will be negative (since $N_{d}>N_{c}$, degenerate semiconductor), indicating that the Fermi level will be somewhere in the conduction band. Actually, the upper expression is derived on the basis of the assumption that the Fermi level is several $k T$ below the bottom of the conduction band, which allows us to replace the Fermi-Dirac statistics with the more simple Boltzmann statistics, so it is not applicable to degenerate semiconductors.

All we can say about the position of the Fermi level for this case is that $E_{F n}>E_{c}-k T$. We can very roughly take that $n=N_{d}$ and estimate the conductivity and the resistivity of the sample,

$$
\begin{gathered}
\sigma \approx e N_{d} \mu_{e}=\left(1.608 \times 10^{-19} \mathrm{C}\right)\left(1 \times 10^{18} \mathrm{~cm}^{-3}\right)\left(2400 \mathrm{~cm}^{2} \mathrm{~V}^{-1} \mathrm{~s}^{-1}\right) \\
=384.5 \Omega^{-1} \mathrm{~cm}^{-1}=3.845 \times 10^{4} \Omega^{-1} \mathrm{~m}^{-1}
\end{gathered}
$$

$$
\rho=\frac{1}{\sigma}=\frac{1}{\left(384.5 \Omega^{-1} \mathrm{~cm}^{-1}\right)}=2.6 \times 10^{-3} \Omega \mathrm{~cm}=2.6 \times 10^{-5} \Omega \mathrm{~m}
$$

NOTE
The exact position of the Fermi level and the exact electron concentration can be determined more rigorously applying the Fermi-Dirac statistics. The semiconductor is described by the following equations:

Electron concentration

$$
n=N_{c} F_{1 / 2}\left(\frac{E_{F n}-E_{c}}{k_{B} T}\right)
$$

Hole concentration (Boltzmann statistic, because $E_{F n}>E_{v}$ )

$$
p=N_{v} \exp \left(\frac{E_{v}-E_{F n}}{k_{B} T}\right)
$$

Concentration of the ionized donors

$$
N_{d}^{+}=\frac{N_{d}}{1+g_{d} \exp \left(\frac{E_{F n}-E_{d}}{k_{B} T}\right)}
$$

No acceptors

$$
N_{a}^{-}=N_{a}=0
$$

Charge neutrality

$$
N_{d}^{+}+p-n=0
$$

where

$$
F_{1 / 2}(\eta)=\frac{2}{\sqrt{\pi}} \int_{0}^{\infty} \frac{x^{1 / 2}}{1+\exp (x-\eta)} d x
$$

is the Fermi-Dirac integral of order $1 / 2$ and $g_{d}$ is the donor degeneracy factor.
Assuming donor ionization energy of 6 meV and degeneracy factor of 2 we can determine the position of the Fermi level and the electron concentration by solving the following equation:

$$
\frac{N_{d}}{1+g_{d} \exp \left(\frac{E_{F n}-E_{d}}{k_{B} T}\right)}+N_{v} \exp \left(\frac{E_{v}-E_{F n}}{k_{B} T}\right)-N_{c} F_{1 / 2}\left(\frac{E_{F n}-E_{c}}{k_{B} T}\right)=0
$$

The Fermi level position and the corresponding electron concentration are presented in the figure below.

$$
\sigma=\text { en } \mu_{e}=\left(1.608 \times 10^{-19} \mathrm{C}\right)\left(3.223 \times 10^{17} \mathrm{~cm}^{-3}\right)\left(2400 \mathrm{~cm}^{2} \mathrm{~V}^{-1} \mathrm{~s}^{-1}\right)
$$

$$
=124 \Omega^{-1} \mathrm{~cm}^{-1}=1.24 \times 10^{4} \Omega^{-1} \mathrm{~m}^{-1}
$$

$$
\rho=\frac{1}{\sigma}=\frac{1}{\left(123.9 \Omega^{-1} \mathrm{~cm}^{-1}\right)}=8.07 \times 10^{-3} \Omega \mathrm{~cm}=8.07 \times 10^{-5} \Omega \mathrm{~m}
$$

These values are a factor of 2-3 different than the rough estimates calculated above.
3.5 AIGaAs LED emitter An AlGaAs LED emitter for use in a local optical fiber network has the output spectrum shown in Figure 3Q5. It is designed for peak emission at 820 nm at $25^{\circ} \mathrm{C}$.
a What is the linewidth $\Delta \lambda$ between half power points at temperatures $-40^{\circ} \mathrm{C}, 25^{\circ} \mathrm{C}$ and $85^{\circ} \mathrm{C}$ ? What is the empirical relationship between $\Delta \lambda$ and $T$ given three temperatures and how does this compare with $\Delta(h v) \approx 2.5 k_{B} T-3 k_{B} T$ ?
b Why does the peak emission wavelength increase with temperature?
c Why does the peak intensity decrease with temperature?
d What is the bandgap of AlGaAs in this LED?
e The bandgap, $E_{g}$, of the ternary alloys $\mathrm{Al}_{x} \mathrm{Ga}_{1-x} \mathrm{As}$ follows the empirical expression,

$$
E_{g}(e V)=1.424+1.266 x+0.266 x^{2}
$$

What is the composition of the AlGaAs in this LED?
f When the forward current is 40 mA , the voltage across the LED is 1.5 V and the optical power that is coupled into a multimode fiber through a lens is $25 \mu \mathrm{~W}$. What is the efficiency?

Relative spectral output power


The output spectrum from GaAlAs LED. Values normalized to peak emission at $25^{\circ} \mathrm{C}$.

Wavelength ( nm )

## Figure 3Q5

## Solution

a We note that the emitted wavelength is related to the photon energy $E_{p h}$ by

$$
\lambda=c / v=h c / E_{p h} .
$$

If we differentiate $\lambda$ with respect to photon energy $E_{p h}$ we get

$$
\frac{d \lambda}{d E_{p h}}=-\frac{h c}{E_{p h}^{2}}
$$

We can represent small changes or intervals (or $\Delta$ ) by differentials, e.g. $\Delta \lambda \Delta E_{p h} \approx$ $\left|d \lambda d E_{p h}\right|$, then

$$
\Delta \lambda \approx \frac{h c}{E_{p h}^{2}} \Delta E_{p h}
$$

We are given the energy width of the output spectrum, $\Delta E_{p h}=\Delta(h v) \approx 3 k_{B} T$. Then, using the latter and substituting for $E_{p h}$ in terms of $\lambda$ we find,

$$
\Delta \lambda \approx \lambda^{2} \frac{3 k_{B} T}{h c} \text { or } \Delta \lambda \approx \lambda^{2} \frac{3 k_{B} T}{h c}
$$

| Temperature $\rightarrow$ <br> Parameter $\downarrow$ | $-40^{\circ} \mathrm{C}$ | $25^{\circ} \mathrm{C}$ | $85^{\circ} \mathrm{C}$ | Comment |
| :--- | :--- | :--- | :--- | :--- |
| $\lambda_{\text {peak }}(\mathrm{nm})$ | 804 | 820 | 837 |  |
| $\Delta \lambda(\mathrm{~nm})$ (Measured) | 30 | 40 | 48 |  |
| $\Delta \lambda(\mathrm{~nm})$ (Calculated); $\Delta E_{p h}=2.5 \mathrm{kT}$ | 26.2 | 34.8 | 43.6 |  |
| $\Delta \lambda(\mathrm{~nm})\left(\right.$ Calculated); $\Delta E_{p h}=3 \mathrm{kT}$ | 31.4 | 41.7 | 52.3 | Very close |



The theory predicts that $\Delta \lambda / \lambda^{2}$ vs $T$ should be a straight line because,

$$
\Delta \lambda \approx \lambda^{2} \frac{\Delta E_{p h}}{h c}
$$

so that $\quad \frac{\Delta \lambda}{\lambda^{2}}=\frac{m k_{B} T}{h c}$
where $\Delta E_{p h}=m k_{B} T$ and $m$ is a numerical constant that represents the ratio $\Delta(h v) /\left(k_{B} T\right)$ and is determined from the slope of the $\Delta \lambda / \lambda^{2}$ vs. $T$ plot.

The three points plotted in the figure seems to follow this behavior. The best line forced through zero has a slope that indicates $m=2.8$.
b The bandgap decreases with temperature
c There are two factors to consider. (i) Spectral intensity means intensity per unit wavelength, that is, $d / / d \lambda$. The integration of the spectral curve gives the total intensity, the total number of photons emitted per unit area per unit time. As the spectrum broadens with temperature we would naturally expect the peak to decrease with temperature. (ii) Higher the temperature, the stronger are the lattice vibrations (there are more phonons). Indirect or radiationless transitions, those that do not emit photons, require phonons (lattice vibrations) which encourage indirect transitions. Thus increasing the temperature increases indirect transitions at the expense of direct transitions and the light intensity decreases. If (ii) was totally absent then the areas under the curves for all the three spectra would be identical.

## d Use the peak emission wavelength to find $E_{g}$ as follows:

$$
\text { At }-40^{\circ} \mathrm{C}(233 \mathrm{~K}), \lambda_{\text {peak }}=804 \mathrm{~nm} .
$$

At $25^{\circ} \mathrm{C}(298 \mathrm{~K}), \lambda_{\text {peak }}=820 \mathrm{~nm}$.
At $85^{\circ} \mathrm{C}(358 \mathrm{~K}), \lambda_{\text {peak }}=837 \mathrm{~nm}$.
We first note that we need the required bandgap $E_{g}$ at the wavelength of interest. The photon energy at peak emission is $h c / \lambda_{\text {peak }}=E_{g}+k_{B} T$. Then,

$$
E_{g}=\frac{c h}{e \lambda_{\text {peak }}}-\frac{k_{B} T}{e}
$$

and at $\lambda_{\text {peak }}=820 \times 10^{9} \mathrm{~m}$, taking $T=25+273 \mathrm{~K}$,

$$
E_{g}=\frac{\left(3 \times 10^{8}\right)\left(6.626 \times 10^{-34}\right)}{\left(1.6 \times 10^{-19}\right)\left(820 \times 10^{-9}\right)}-0.0257 \mathrm{eV}=\mathbf{1 . 4 8 6 3} \mathbf{~ e V}
$$

e The bandgap $E_{g}$ of the ternary alloys $\mathrm{Al}_{x} \mathrm{Ga}_{1-x} \mathrm{As}$ follows the empirical expression,

$$
\begin{aligned}
& E_{g}(\mathrm{eV})=1.424+1.266 x+0.266 x^{2} . \\
\therefore \quad & E_{g}(e V)=1.4863=1.424+1.266 x+0.266 x^{2} .
\end{aligned}
$$

Solving for $x$ we find, $x=\mathbf{0 . 0 5}$.
f From the definition of efficiency $\eta$,

$$
\begin{aligned}
& \eta=\frac{\text { Output optical power }}{\text { Input electrical power }} \\
& \qquad=\frac{P_{o}}{I V}=\frac{25 \times 10^{-6} \mathrm{~W}}{\left(40 \times 10^{-3} \mathrm{~A}\right)(1.5 \mathrm{~V})}=0.000417=0.0417 \%
\end{aligned}
$$

3.6 III-V compound semiconductors in optoelectronics Figure 3Q6 represents the bandgap $E_{g}$ and the lattice parameter $a$ in the quarternary III-V alloy system. A line joining two points represents the changes in $E_{g}$ and $a$ with composition in a ternary alloy composed of the compounds at the ends of that line. For example, starting at GaAs point, $E_{g}=1.42 \mathrm{eV}$ and $a=$ 0.565 nm , and $E_{g}$ decreases and $a$ increases as GaAs is alloyed with InAs and we move along the line joining GaAs to InAs. Eventually at InAs, $E_{g}=0.35 \mathrm{eV}$ and $a=0.606 \mathrm{~nm}$. Point $X$ in Figure 3Q6 is composed of InAs and GaAs and it is the ternary alloy $\mathrm{In}_{x} \mathrm{Ga}_{1-x} \mathrm{As}$. It has $E_{g}=0.7 \mathrm{eV}$ and $a=0.587 \mathrm{~nm}$ which is the same $a$ as that for $\operatorname{InP} \operatorname{In}_{x} \mathrm{Ga}_{1-x} \mathrm{As}$ at $X$ is therefore lattice matched to InP and can hence be grown on an InP substrate without creating defects at the interface.

Further, $\mathrm{In}_{x} \mathrm{Ga}_{1-x} \mathrm{As}$ at $X$ can be alloyed with InP to obtain a quarternary alloy $\mathrm{In}_{x} \mathrm{Ga}_{1-}$ ${ }_{x} \mathrm{As}_{y} \mathrm{P}_{1-y}$ whose properties lie on the line joining $X$ and InP and therefore all have the same lattice parameter as $\operatorname{InP}$ but different bandgap. Layers of $\operatorname{In}_{x} \mathrm{Ga}_{1-x} \mathrm{As}_{y} \mathrm{P}_{1-y}$ with composition between $X$ and InP can be grown epitaxially on an InP substrate by various techniques such as liquid phase epitaxy (LPE) or molecular beam expitaxy (MBE).

The shaded area between the solid lines represents the possible values of $E_{g}$ and $a$ for the quarternary III-V alloy system in which the bandgap is direct and hence suitable for direct recombination.

The compositions of the quarternary alloy lattice matched to $\operatorname{InP}$ follow the line from $X$ to InP.
a Given that the $\operatorname{In}_{x} \mathrm{Ga}_{1-x}$ As at $X$ is $\operatorname{In}_{0.535} \mathrm{Ga}_{0.465}$ As show that quarternary alloys $\operatorname{In}_{x} \mathrm{Ga}_{1-}$ ${ }_{x} \mathrm{As}_{y} \mathrm{P}_{1-y}$ are lattice matched to $\operatorname{InP}$ when $y=2.15 x$.
b The bandgap energy $E_{g}$, in eV for $\mathrm{In}_{x} \mathrm{Ga}_{1-x} \mathrm{As}_{y} \mathrm{P}_{1-y}$ lattice matched to InP is given by the empirical relation,

$$
E_{g}(e V)=1.35-0.72 y+0.12 y^{2}
$$

Find the composition of the quarternary alloy suitable for an emitter operating at 1.55 $\mu \mathrm{m}$.


Bandgap energy $E_{g}$ and lattice constant $a$ for various III-V alloys of GaP, GaAs, InP and InAs. A line represents a ternary alloy formed with compounds from the end points of the line. Solid lines are for direct bandgap alloys whereas dashed lines for indirect bandgap alloys. Regions between lines represent quaternary alloys. The line from $X$ to InP represents quaternary alloys $\mathrm{In}_{1-x} \mathrm{Ga}_{x} \mathrm{As}_{1-y} \mathrm{P}_{y}$ made from $\mathrm{In}_{0.535} \mathrm{Ga}_{0.465} \mathrm{As}$ and InP which are lattice matched to InP.

Figure 3Q6

Solution
a The III-V quaternary alloy is made of $\left(\operatorname{In}_{0.535} \mathrm{Ga}_{0.465} \mathrm{As}\right)$ alloyed with $\operatorname{InP}$. Suppose that $z$ fraction of $\mathrm{In}_{0.535} \mathrm{Ga}_{0.465} \mathrm{As}$ is alloyed with $(1-z)$ fraction of $\operatorname{InP}$.
$z=0$ for $\operatorname{InP}$ and $z=1$ for $\operatorname{In}_{0.535} \mathrm{Ga}_{0.465} \mathrm{As}$
The formula for the quaternary alloy is $\mathrm{In}_{1-x} \mathrm{Ga}_{x} \mathrm{As}_{y} \mathrm{P}_{1-y}$.
Thus,

$$
z\left(\mathrm{In}_{0.535} \mathrm{Ga}_{0.465} \mathrm{As}\right)+(1-z) \mathrm{InP}=\mathrm{In}_{1-x} \mathrm{Ga}_{x} \mathrm{As}_{y} \mathrm{P}_{1-y}
$$

Use $x=x^{\prime}$ to represent point X or $\left(\operatorname{In}_{0.535} \mathrm{Ga}_{0.465} \mathrm{As}\right)$, that is, $x^{\prime}=0.535$ for the In content.
We can now balance the contents for each element.
For P

$$
1-z=1-y \quad \therefore \quad z=y
$$

For $\mathrm{Ga} \quad\left(1-x^{\prime}\right) z=x \quad \therefore \quad y=2.15 x$
For In $\quad x^{\prime} z+(1-\mathrm{z})=1-x \quad \therefore \quad z=x /\left(1-x^{\prime}\right)$
For As $\quad z=y \quad \therefore \quad y=2.15 x$
b Consider a lattice matched InGaAsP to InP. Given $\lambda=1.55 \mu \mathrm{~m}$.

$$
E_{g}=1.24 / \lambda=0.80 \mathrm{eV}
$$

Thus,

$$
E_{g}(e V)=1.35-0.72 y+0.12 y^{2}=0.80 \mathrm{eV}
$$

$$
\therefore \quad 0.12 y^{2}-0.72 y+0.55=0
$$

$$
\therefore \quad y=0.90, x=0.42
$$

$$
\therefore \quad \mathrm{In}_{0.58} \mathrm{Ga}_{0.42} \mathrm{As}_{0.9} \mathrm{P}_{0.1}
$$

3.8 Linewidth of LEDs Experiments carried out on various direct bandgap semiconductor LEDs give the output spectral linewidth (between half intensity points) listed in Table 3Q8-1. We know that a spread in the wavelength is related to a spread in the photon energy,

$$
\begin{equation*}
\Delta \lambda \approx \frac{h c}{E_{p h}^{2}} \Delta E_{p h} \tag{1}
\end{equation*}
$$

Suppose that we write $E_{p h}=h c / \lambda$ and $\Delta E_{p h}=\Delta(h v) \approx m k_{B} T$ where $m$ is a numerical constant. Show that,

$$
\begin{equation*}
\Delta \lambda \approx \lambda^{2} \frac{m k_{B} T}{h c} \tag{2}
\end{equation*}
$$

and by appropriately plotting the data in Table 3Q8-1 find $m$.
Table 3Q8-1 shows the linewidth $\Delta \lambda_{1 / 2}$ for various visible LEDs. Radiative recombination is obtained by appropriately doping the material. Using $m \approx 3, T=300 \mathrm{~K}$, in Eq. (2) calculate the expected spectral width for each and compare with the experimental value. What is your conclusion? Do you think $E_{N}$ in Figure 3.24 (b) is a discrete level?

## Table 3Q8-1

Linewidth $\Delta \lambda_{1 / 2}$ between half points in the output spectrum (Intensity vs wavelength) of GaAs and AlGaAs LEDs.

| Peak wavelength of <br> emission $(\lambda) \mathrm{nm}$ | 650 | 810 | 820 | 890 | 950 | 1150 | 1270 | 1500 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\Delta \lambda_{1 / 2} \mathrm{~nm}$ | 22 | 36 | 40 | 50 | 55 | 90 | 110 | 150 |
| Material (Direct $E_{g}$ ) | AlGaAs | AlGaAs | AlGaAs | GaAs | GaAs | InGaAsP | InGaAsP | InGaAsP |

## Table 3Q8-2

Linewidth $\Delta \lambda_{1 / 2}$ between half points in the output spectrum (Intensity vs wavelength) of four various visible LEDs using SiC and GaAsP materials.

| Peak wavelength of emission $(\lambda) \mathrm{nm}$ | 468 | 565 | 583 | 600 | 635 |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $\Delta \lambda_{1 / 2} \mathrm{~nm}$ | 66 | 28 | 36 | 40 | 45 |
| Color | Blue | Green | Yellow | Orange | Red |
| Material | SiC (Al) | GaP (N) | GaAsP (N) | GaAsP (N) | GaAsP |

Solution


The best line forced through zero is

$$
\Delta \lambda_{1 / 2}=\left[6.57 \times 10^{-5}(\mathrm{~nm})^{-1}\right] \lambda^{2}
$$

and

$$
\text { Slope }=6.57 \times 10^{-5}(\mathrm{~nm})^{-1} .
$$

This slope is $m k T / h c$; thus $m \approx 3.15$. Very close to the theoretically predicted value of $m=2.5-$ 3.

Table 2Q8-3 shows the calculated spectral widths $\Delta \lambda_{1 / 2}$ using $m=3$. The actual observed widths are substantially larger than the expected $\Delta \lambda_{1 / 2}$ using $m=3$, as we found above for the direct bandgap materials. If the recombination center were a discrete level, we would expect a spread in the photon energy that is controlled by the energy distribution of holes
(electrons) in the valence (conduction) band, that is about $\sim 1.5 k_{B} T$, or $m=1.5$. The observed spread is much more than $m=1.5$ and hence the energy level cannot be discrete. ( It is possible to give a semiquantitative plausible explanation as follows. A captured electron will have a wavefunction that is localized and hence a smaller uncertainty in its position than in the band; i.e. $\Delta x$ will be small. That means the uncertainty $\Delta p$ in its momentum will be higher and hence the uncertainly in its energy will also be higher. We would expect that the spread of photon energies will be more than from than in band to band recombinations.)

## Table 3Q8-3

Linewidth $\Delta \lambda_{1 / 2}$ between half points in the output spectrum (Intensity vs wavelength) of four various visible LEDs using SiC and GaAsP materials.

| Peak wavelength of emission $(\lambda) \mathrm{nm}$ | 468 | 565 | 583 | 600 | 635 |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $\Delta \lambda_{1 / 2} \mathrm{~nm}$ | 66 | 28 | 36 | 40 | 45 |
| Expected $\Delta \boldsymbol{\lambda}_{\mathbf{1} / \mathbf{2}} \mathbf{n m}$ using $\boldsymbol{m}=\mathbf{3}$ | $\mathbf{1 3 . 7}$ | $\mathbf{2 0 . 0}$ | $\mathbf{2 1 . 3}$ | $\mathbf{2 2 . 5}$ | $\mathbf{2 5 . 2}$ |
| Color | Blue | Green | Yellow | Orange | Red |
| Material | SiC (Al) | GaP (N) | GaAsP (N) | GaAsP (N) | GaAsP |

### 3.10 LED-Fiber coupling Efficiency

a It is found that approximately $200 \mu \mathrm{~W}$ is coupled into a multimode step index fiber from a surface emitting LED when the current is 75 mA and the voltage across the LED is about 1.5 V. What is the overall efficiency of operation?
b Experiments are carried out on coupling light from a 1310 nm ELED (edge emitting LED) in multimode and single mode fibers.
(i) At room temperature, when the ELED current is 120 mA , the voltage is 1.3 V and light power coupled into a $50 \mu \mathrm{~m}$ multimode fiber with NA (numerical aperture) $=0.2$ is $48 \mu \mathrm{~W}$. What is the overall efficiency?
(ii) At room temperature, when the ELED current is 120 mA , the voltage is 1.3 V and light power coupled into a $9 \mu \mathrm{~m}$ single mode fiber is $7 \mu \mathrm{~W}$. What is the overall efficiency?

## Solution

a

$$
\eta_{\text {overall }}=\frac{P_{o}}{I V}=\frac{200 \times 10^{-6} \mathrm{~W}}{\left(75 \times 10^{-3} \mathrm{~A}\right)(1.5 \mathrm{~V})}=1.8 \times 10^{-3}=\mathbf{0 . 1 8} \%
$$

b (i)

$$
\eta_{\text {overall }}=\frac{P_{o}}{I V}=\frac{48 \times 10^{-6} \mathrm{~W}}{\left(120 \times 10^{-3} \mathrm{~A}\right)(1.3 \mathrm{~V})}=\mathbf{0 . 0 3 0 7 \%}
$$

(ii) $\quad \eta_{\text {overall }}=\frac{P_{o}}{I V}=\frac{7 \times 10^{-6} \mathrm{~W}}{\left(120 \times 10^{-3} \mathrm{~A}\right)(1.3 \mathrm{~V})}=\mathbf{0 . 0 0 4 5 \%}$

